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A high temperature phase of  $(\text{Ni}_{1-x}\text{Mn}_x)_{1.5}\text{Ge}$  is found to be a hexagonal  $\text{Ni}_2\text{In}$  type structure in the range of  $0.5 \leq x \leq 0.7$ . A phase transition of  $\text{Ni}_2\text{In}$  type  $\Leftrightarrow$   $\text{TiNiSi}$  type occurs in the alloys of  $0.45 \leq x \leq 0.55$  below 270K.  $\text{Ni}_{1.5}\text{Ge}$  is paramagnetic and a magnetic correlation is induced with substitution of Mn for Ni. In the alloys of  $0.5 \leq x < 0.35$ , the temperature dependence of ac susceptibility ( $\chi''$ - $T$  curve) shows a cusp first and then a hump with substitution of Mn for Ni. The cusp and hump suggest a change from a spin glass state and a cluster spin glass state with the substitution. The alloys with  $0.35 \leq x < 0.45$  are ferromagnetic. The  $\sigma$ - $T$  curve on cooling shows ferromagnetism at  $T_C$  and then a magnetic cooling effect below the freezing temperature  $T_f$ , suggesting a re-entrant spin glass state. The glass like behavior shows an existence of antiferromagnetic interactions between the nearest neighbors in the hexagonal lattice.  $(\text{NiMn})_{1.5}\text{Ge}$  ( $x=0.5$ ) exhibits the  $\text{Ni}_2\text{In}$  type  $\Leftrightarrow$   $\text{TiNiSi}$  type phase transition at  $T_D=267\text{K}$  on cooling, while the stoichiometric  $\text{NiMnGe}$  shows the phase transition at  $T=470\text{K}$ . The  $\text{Ni}_2\text{In}$  type structure is stabilized with the deficiency of transition metal atoms. The  $\text{TiNiSi}$  type phase is an antiferromagnet, perhaps of a spiral magnetic structure as same as  $\text{NiMnGe}$ . The alloys of  $x=0.45$  and  $0.55$  show paramagnetism, ferromagnetism and antiferromagnetism on cooling and this process is reversible. The ferromagnetism  $\Leftrightarrow$  antiferromagnetism transition occurs with the  $\text{Ni}_2\text{In}$   $\Leftrightarrow$   $\text{TiNiSi}$  phase transition. The alloys of  $0.55 \leq x \leq 0.7$  are the  $\text{Ni}_2\text{In}$  type structure and show ferromagnetism. The Curie temperature decreases from  $T_C=204\text{K}$  of  $x=0.55$  to  $148\text{K}$  of  $x=0.7$ .

#### Q2-60 METAMAGNETISM IN Fe-Mn-Al ALLOYS

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In ternary alloys consisting of two 3d transition elements (i.e. Fe, Mn) and aluminium (Al) admixtures of the last element cause drastic changes in the magnetic moments of the constituents and therefore on the alloy magnetic properties. In this contribution we report results of experiments performed on a series of  $\text{Fe}_{1-x}\text{Mn}_x\text{Al}_x$  alloys by magnetisation measurements and Mössbauer spectroscopy. Increasing step by step the aluminium concentration the alloys change their behaviour from ferromagnetic to antiferromagnetic crossing regions with re-entrant spin glass behaviour. In regions of aluminium concentrations above 25 at% we observe the onset of antiferromagnetism at temperatures around 20 K. Application of magnetic fields leads to typical metamagnetic behaviour. Increasing the magnetic field causes the cusp (characteristic for antiferromagnetism) disappearing and the magnetisation value at low temperatures increasing. An important experimental fact is that these alloys exhibit some tendencies to order. Because of the coexistence of a disordered and ordered phase the interesting experimental question is which phase is causing the metamagnetic behaviour. Additionally Mössbauer spectroscopy is applied without external magnetic fields. Observing the spectra below the onset of antiferromagnetism versus temperature unexpected courses in the isomer shift and the spectrum area appear. Both quantities indicate changes in the electronic structure of these alloys due to the magnetic phase transition.

#### Q2-61 MAGNETIC HYPERFINE FIELD AT Hf SITE IN $\text{Hf}(\text{Fe}_{1-x}\text{Co}_x)_2$ AND $\text{Hf}(\text{Fe}_{1-x}\text{Co}_x)_2\text{H}_y$ AT LOW Co CONCENTRATION.

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The stoichiometric  $\text{AB}_2$  cubic Laves compounds are well suited for the study of the magnetic hyperfine field behaviour because of their cubic symmetry and well defined position of the atoms in the lattice. The pseudobinary compound  $\text{Hf}(\text{Fe}_{1-x}\text{Co}_x)_2$  crystallizes in the cubic Laves phase showing a spontaneous magnetization which decreases with  $x$  and presents a transition from ferromagnetism to paramagnetism at  $x = 0.65$ , with a sharp collapse of the magnetic moment[1].

At low Co concentration, different Hf nearest neighbour configurations can simultaneously be present in several sublattices varying from undisturbed neighbourhood (only Fe atoms) to configurations where one or more Co replace Fe atoms. The measurement of the magnetic hyperfine field at the Hf site in these situations is therefore an ideal way to investigate microscopically the collapse of the ferromagnetic order.

In the present work, the magnetic hyperfine field has been investigated at Hf site in pseudobinary  $\text{Hf}(\text{Fe}_{1-x}\text{Co}_x)_2$  and  $\text{Hf}(\text{Fe}_{1-x}\text{Co}_x)_2\text{H}_y$  ( $0.1 \leq y < 1$ ) alloys in the range of  $0 \leq x \leq 0.6$  by TDPAC measurements. The results show that for  $0.01 \leq x \leq 0.10$  there are three different values of mhf, corresponding to undisturbed neighbourhood of the Hf, and the Hf neighbourhood configurations where one

and two or more Co atoms replace some of the 12 nearest neighbour Fe atoms respectively. The value of the mhf at Hf site for undisturbed neighbourhood decreases linearly with the Co concentration, which indicates that a longer range rather than just the nearest neighbour interaction is important in the formation of the magnetic order. The results for hydrogenated alloys show a similar behaviour except that the presence of long range interactions is not indicated. [1]T. Sakakibara, T. Goto, Y. Nihuhara, J. de Phys. 12(1988)C8-263

#### Q2-62 COMPARISON OF THE MOSSBAUER SPECTRA OF BULK $\text{Fe}_{0.5x}\text{Ni}_{0.4x}$ TO THOSE OF ULTRAFINE $\text{Fe}_{0.50}\text{Ni}_{0.50}$ .

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Mössbauer spectra of bulk  $\text{Fe}_{0.5x}\text{Ni}_{0.4x}$  and 10 nm  $\text{Fe}_{0.50}\text{Ni}_{0.50}$  were collected in zero external magnetic field at temperatures between 2.6 K and room temperature, and at 6 K in longitudinal external magnetic fields between 1.3 T and 14.3 T. The two sets of spectra were compared in order to determine the surface attributes of ultrafine iron-nickel grains. The spectra of 10 nm  $\text{Fe}_{0.50}\text{Ni}_{0.50}$  have been interpreted with reference to the characteristics ascribed to nanocrystalline matter.

#### Q2-63 TRANSVERSE SPIN FREEZING IN ANTIFERROMAGNETIC $\gamma$ - $\text{Mn}_2\text{Cu}_{10}$

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$\gamma$ - $\text{Mn}_2\text{Cu}_{10}$  is face centred tetragonal with  $c/a < 1$  and antiferromagnetic with the antiferromagnetic direction along the  $c$ -direction. A single domain can be preferred in a single crystal by applying pressure along the  $c$ -direction. A crystal prepared in this way was checked by neutron diffraction and found to be more than 98% one domain. Low field (100 Oe) field cooled (FC) and zero-field cooled (ZFC) magnetic susceptibility measurements clearly show spin freezing at 50K, well below the Néel temperature of about 425K. Measurements taken parallel to the antiferromagnetic direction show much less difference between the ZFC and FC susceptibilities than those taken perpendicular to the antiferromagnetic direction along  $\langle 110 \rangle$ . Preliminary neutron measurements of the temperature dependence of the intensity of the (110) magnetic Bragg peak show no anomaly at the freezing temperature.

#### Q2-64 THE TEMPERATURE DEPENDENCE OF THE MAGNETIC EXCITATIONS IN NEARLY ANTIFERROMAGNETIC $\text{Cr}_{0.95}\text{V}_{0.05}$

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$\text{Cr}_{0.95}\text{V}_{0.05}$  is close to antiferromagnetic order at low temperatures. We present neutron scattering measurements of the magnetic excitations in this material using triple-axis and pulsed-spallation-neutron scattering techniques. Measurements are made in the energy and temperature ranges 2-400 meV and 2-1000 K, respectively. In agreement with previous measurements, we observe six incommensurate peaks in  $c^*(Q,w)$  at low frequency and temperature. These have been previously attributed to Fermi surface nesting. At all temperatures, the low-frequency wavevector-integrated susceptibility  $c^*(w)$  is proportional to  $w$ . On warming, the incommensurate peaks merge into a single commensurate peak and the gradient of  $c^*(w)$  is substantially reduced. Increasing frequency also results in a broadening of the peaks in  $c^*(Q,w)$ . Comparisons are made with measurements on the high- $T_C$  superconductor  $\text{La}_{1.86}\text{Sr}_{0.14}\text{CuO}_4$  where quantitatively similar observations have been made.

#### Q2-65 ATOMIC-VOLUME DEPENDENCE OF SPIN STATE IN Fe-RICH FCC Fe-Pt-Ru ALLOYS

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